Integral Representations in Weighted Bergman Spaces on the Tube Domains

Yun Huang ^{*}, Guan-Tie Deng [†], Tao Qian [‡]

Herein, the Laplace transform representations for functions of weighted holomorphic Bergman spaces on the tube domains are developed. Then a weighted version of the edge-of-the-wedge theorem is derived as a byproduct of the main results.

Key words: Weighted Bergman space, Tube domain, Laplace transform, Integral representation, Regular cone

1 Introduction

The classical Paley–Wiener theorem asserts that functions of the classical Hardy space $H^2(\mathbb{C}^+)$ can be written as the Laplace transforms of L^2 functions supported in the right half of the real axis, see [1]. This theorem has been extended to more general Hardy spaces, including the H^p spaces cases (0), higher dimensional cases and weighted spaces, see [11, 13, 15, 14, 12, 9]. Integral representation theorems have also been investigated for Bergman spaces.

^{*}Department of Mathematics, Faculty of Science and Technology, University of Macau, Macao (Via Hong Kong). Email: yhuang12@126.com.

[†]Corresponding author. School of Mathematical Sciences, Beijing Normal University, Beijing, 100875. Email:denggt@bnu.edu.cn. This work was partially supported by NSFC (Grant 11971042) and by SRFDP (Grant 20100003110004)

[‡]Institute of Systems Engineering, Macau University of Science and Technology, Avenida Wai Long, Taipa, Macau). Email: tqian1958@gmail.com. The work is supported by the Macau Science and Technology foundation No.FDCT079/2016/A2, FDCT0123/2018/A3, and the Multi-Year Research Grants of the University of Macau No. MYRG2018-00168-FST.

We first introduce some notations and definitions. Let B be a domain (open and connected set) in \mathbb{R}^n and $T_B = \mathbb{R}^n + iB \subset \mathbb{C}^n$ be the tube over B. For any element $z = (z_1, z_2, \ldots, z_n), z_k = x_k + iy_k$, by definition, $z \in T_B$ is and only if $x = (x_1, \ldots, x_n) \in \mathbb{R}^n$ and $y = (y_1, \ldots, y_n) \in B$. The inner product of $z, w \in \mathbb{C}^n$ is defined as $z \cdot w = z_1 w_1 + z_2 w_2 + \cdots + z_n w_n$. The associated Euclidean norm of z is $|z| = \sqrt{z \cdot \overline{z}}$, where $\overline{z} = (\overline{z}_1, \overline{z}_2, \ldots, \overline{z}_n)$.

A nonempty subset $\Gamma \subset \mathbb{R}^n$ is called an open cone if it satisfies (i) $0 \notin \Gamma$, and (ii) $\alpha x + \beta y \in \Gamma$ for any $x, y \in \Gamma$ and $\alpha, \beta > 0$. The dual cone of Γ is defined as $\Gamma^* = \{y \in \mathbb{R}^n : y \cdot x \ge 0, \text{ for any } x \in \Gamma\}$, which is clearly a closed convex cone with vertex at 0. We say that the cone Γ is regular if the interior of its dual cone Γ^* is nonempty.

For $\frac{1}{p} + \frac{1}{q} = 1$, define

$$B^{p}(T_{B}) = \left\{ F : F \text{ is holomorphic in } T_{B} \text{ and satisfies } \int_{B} \left(\int_{\mathbb{R}^{n}} |F(x+iy)|^{p} dx \right)^{q-1} dy < \infty \right\}.$$

Among the previous studies, Genchev showed that the function spaces $B^p(1 \le p \le 2)$, in the one- and multi-dimensions in [3] and [4], respectively, admit integral representations in the Laplace transform form. These results can be applied to the Bergman spaces

$$A^{p}(T_{\Gamma}) = \left\{ F : F \text{ is holomorphic on } T_{\Gamma} \text{ and satisfies } \int_{T_{\Gamma}} |F(x+iy)|^{p} dx dy < \infty \right\}$$

to obtain the corresponding integral representation results for $A^p(T_{\Gamma})$ in the range $1 \le p \le 2(5)$.

In this paper we initiate a study on a class of function spaces, denoted by $A^{p,s}(B,\psi)$, of which each is associated with a weight function of the form $e^{-2\pi\psi(y)}$, where $\psi(y) \in C(B)$ is continuous on B. The space $A^{p,s}(B,\psi)(0 is the collection of functions$ <math>F(z) that are holomorphic in T_B and satisfy

$$\|F\|_{A^{p,s}(B,\psi)} = \left(\int_B \left(\int_{\mathbb{R}^n} |F(x+iy)e^{-2\pi\psi(y)}|^p dx \right)^s dy \right)^{\frac{1}{sp}} < \infty, \ 0 < p, s < \infty,$$
$$|F\|_{A^{p,\infty}(B,\psi)} = \sup\left\{ e^{-2\pi\psi(y)} \left(\int_{\mathbb{R}^n} |F(x+iy)|^p dx \right)^{\frac{1}{p}}, y \in B \right\} < \infty, \ 0 < p < \infty, s = \infty$$

and

$$||F||_{A^{\infty,\infty}(B,\psi)} = \sup\left\{e^{-2\pi\psi(y)}|F(x+iy)|, x \in \mathbb{R}^n, y \in B\right\} < \infty, \ p = \infty, s = \infty.$$

This paper is structured as follows. In §2, we introduce our main work on the integral representation for $A^{p,s}(B,\psi)$, which is separated into three cases, namely, $A^{p,s}(B,\psi)$ for

 $1 \le p \le 2$, $A^{p,s}(B,\psi)$ for $0 and <math>A^{p,s}(\Gamma,\psi)$ for p > 2, corresponding to Theorem 1, 2 and 3 respectively. The proofs are given in §3. Finally, some results, referring to Corollary 2, Theorem 4 and Theorem 5, are derived as applications of the integral representation theorems claimed in §2.

2 Main results

In order to introduce our main results, we define the set

$$U_{\alpha}(B,\psi) = \left\{ t \in \mathbb{R}^{n} : \int_{B} e^{-2\pi\alpha(t \cdot y + \psi(y))} dy < \infty \right\}$$
(1)

for $\alpha \in (0, \infty)$ and

$$U_{\infty}(B,\psi) = \{t : \inf_{y \in \Gamma} (y \cdot t + \psi(y)) > -\infty\}$$
(2)

for $\alpha = \infty$.

The representation theorem for $A^{p,s}(B,\psi)$, where $1 \le p \le 2$ and $0 < s \le \infty$, is stated as follows.

THEOREM 1. Assume that $1 \le p \le 2$, $0 < s \le \infty$, then each $F(z) \in A^{p,s}(B,\psi)$ admits an integral representation in the form

$$F(z) = \int_{\mathbb{R}^n} f(t) e^{2\pi i t \cdot z} dt, \ z \in T_B,$$
(3)

in which, for p = 1, $f(t) \in C(\mathbb{R}^n)$ satisfies

$$|f(t)| \left(\int_{B} e^{-2s\pi(y \cdot t + \psi(y))} dy \right)^{\frac{1}{s}} \le ||F||_{A^{1,s}(B,\psi)}$$
(4)

and, for $1 and <math>\frac{1}{p} + \frac{1}{q} = 1$, f(t) is a measurable function that satisfies

$$\left(\int_{B} \left(\int_{\mathbb{R}^{n}} |f(t)e^{-2\pi(y\cdot t+\psi(y))}|^{q} dt\right)^{\frac{sp}{q}} dy\right)^{\frac{1}{sp}} \leq \|F\|_{A^{p,s}(B,\psi)}.$$
(5)

Moreover, f is supported in $U_s(B, \psi)$ for p = 1 and supported in $U_{sp}(B, \psi)$ for 1 , $<math>0 < s(p-1) \le 1$.

As given in the next theorem, integral representations in the form of Laplace transform are also available for $0 and <math>0 < s \le \infty$.

THEOREM 2. Assume that $F(z) \in A^{p,s}(B, \psi)$, where $0 and <math>0 < s \le \infty$. Then there exists a continuous function f(t) such that $f(t)e^{-2\pi y \cdot t} \in L^1(\mathbb{R}^n)$ and (3) hold for $y \in B$.

Considering the property of f(t) for the case of 0 , we let B be a regular open $convex cone <math>\Gamma$ and let $\psi \in C(\Gamma)$ satisfy

$$R_{\psi} = \lim_{y \in \Gamma, y \to \infty} \frac{\psi(y)}{|y|} < \infty.$$
(6)

Then we obtain the following corollary.

COROLLARY 1. Assume that Γ is a regular open convex cone and $F(z) \in A^{p,s}(\Gamma, \psi)$ for $0 , where <math>\psi \in C(\Gamma)$ satisfies (6). Then there exists f(t) supported in $\Gamma^* + \overline{D(0, R_{\psi})}$ such that (3) holds and $|f(t)| \left(\int_{\Gamma} e^{-2s\pi(y\cdot t + R_{\psi}|y|)} dy\right)^{\frac{1}{s}}$ is slowly increasing.

Similarly, we establish an analogy for p > 2 and $0 < s \le \infty$.

THEOREM 3. Assume that $p > 2, 0 < s \le \infty$, Γ is a regular open convex cone in \mathbb{R}^n and $\psi \in C(\Gamma)$ satisfies (6). If $F(z) \in A^{p,s}(\Gamma, \psi)$ satisfying

$$\lim_{y \in \overline{\Gamma}, y \to 0} \int_{\mathbb{R}^n} |F(x+iy)|^2 dx < \infty, \tag{7}$$

then there exists $f(t) \in L^2(\mathbb{R}^n)$ supported in $U_{sp}(\Gamma, \psi)$ such that (3) holds for all $z \in T_{\Gamma}$.

The definition of $A^{p,s}(B, \psi)$ shows that $A^{p,s}(B, \psi)$ is a weighted Hardy space when $s = \infty$ and a weighted Bergman space when s = 1. Taking $\psi(y) = 0$, it becomes, for $s = \infty$ and s = 1, respectively, the classical Hardy space H^p and the classical Bergman space A^p . Therefore, our results herein can be regarded as generalizations of certain previously obtained results.

For example, taking $s = \infty$ and B a regular open convex cone Γ , $A^{p,\infty}(B,\psi) = H^p(\Gamma,\psi)$ is the weighted Hardy spaces investigated in our previous paper [15]. Then Theorem 1, 2 and 3 in [15] can be derived from our main work, including Theorem 1, 2, 3 and Corollary 1 herein. For $s = \infty$ and $\psi(y) = 0$, letting B be some specific domains, some previous studies for the Hardy spaces, see [1, 13, 14, 12, 9], can be also derived from Theorem 1, 2, 3 and Corollary 1.

On the other hand, letting s = 1, by using Theorem 1, 2, 3 and Corollary 1, we can obtain the representation theorems for the standard Bergman spaces. Note that for s = 1, $B = \Gamma$ and $\psi(y) = 0$, we have $A^{p,s}(B, \psi) = A^p(T_{\Gamma})$. We therefore conclude from Theorem 1 that the counterpart results of Theorem 1, 2 and 3 in [5] hold for the classical Bergman spaces $A^p(T_{\Gamma})(1 \le p \le 2)$. If we set $\psi(y) = 0$ and s = q - 1, where $\frac{1}{p} + \frac{1}{q} = 1$, then $A^{p,s}(B, \psi) = B^p(T_B)$. The integral representation theorems for those function spaces $B^p(T_B)(1 \le p \le 2)$ can be derived from Theorem 1 herein, see [4]. Especially, letting s = 1, p = 2, $\psi(y) = -\frac{\alpha}{4\pi} \log |y|$ and B a regular open convex cone Γ , Theorem 1 implies a higher dimensional generalization of Theorem 1 of [10] in tube domains, which is established as Corollary 2 in the sequel.

3 Proofs

This section is devoted to proving the results stated in §2.

Proof of Theorem 1. We first prove the case of p = 1. If $F(z) \in A^{1,s}(B, \psi)$, then $F_y(x) \in L^1(\mathbb{R}^n)$ as a function of x, and $\check{F}_y(x)$ as well, are both well defined. Next we prove that $\check{F}_y(t)e^{-2\pi y \cdot t}$ is independent of $y \in B$. Without loss of generality, assume that $a = (a', a_n)$, $b = (a', b_n) \in B$, and $a + \tau(b - a) \in B$ for $0 \le \tau \le 1$, where $a' = (a_1, \ldots, a_{n-1})$. The fact $F_y(x) \in L^1(\mathbb{R}^n)$ implies that

$$\int_0^\infty \int_0^1 \int_{\mathbb{R}^{n-1}} \left(|F((x', x_n) + i(a + \tau(b - a)))| + |F((x', -x_n) + i(a + \tau(b - a)))| \right) dx' d\tau dx_n < \infty,$$

which implies

$$\lim_{R \to \infty} \int_0^1 \int_{\mathbb{R}^{n-1}} \left(|F((x', R) + i(a + \tau(b - a)))| + |F((x', -R) + i(a + \tau(b - a)))| \right) dx' d\tau = 0.$$

Hence, we have

$$\begin{split} &|\check{F}_{b}(t)e^{-2\pi b\cdot t} - \check{F}_{a}(t)e^{-2\pi a\cdot t}| \\ &= \left| \int_{\mathbb{R}^{n}} \left(F(x+ib)e^{2\pi i(x+ib)\cdot t} - F(x+ia)e^{2\pi i(x+ia)\cdot t} \right) dx \right| \\ &= \left| \int_{\mathbb{R}^{n}} \int_{0}^{1} \frac{\partial}{\partial \tau} \left(F(x+i(a+\tau(b-a)))e^{2\pi i(x+i(a+\tau(b-a)))\cdot t} \right) d\tau dx \right| \\ &= \left| \int_{\mathbb{R}^{n}} \int_{0}^{1} \frac{\partial}{\partial y_{n}} \left(F(x+i((y',y_{n}))e^{2\pi i(x+i(y',y_{n}))\cdot t} \Big|_{y_{n}=a_{n}+\tau(b_{n}-a_{n})}(b_{n}-a_{n}) \right) d\tau dx \right| \\ &= \left| b_{n} - a_{n} \right| \left| \int_{\mathbb{R}^{n}} \int_{0}^{1} i \frac{\partial}{\partial x_{n}} \left(F(x+i(a+\tau(b-a)))e^{2\pi i(x+i(a+\tau(b-a)))\cdot t} \right) d\tau dx \right| \\ &\leq \left| b_{n} - a_{n} \right| \lim_{R \to \infty} \int_{0}^{1} \int_{\mathbb{R}^{n-1}} \left(\left| F((x',R),(a+\tau(b-a))) \right| + \left| F((x',-R),(a+\tau(b-a))) \right| \right) \right| \\ &e^{-2\pi |t| (|a|+|b-a|)} dx' d\tau \end{split}$$

= 0.

Remark that *B* is connected and open, by an iteration argument, we can show that $\check{F}_y(t)e^{-2\pi y \cdot t}$ is independent of $y \in B$ and we write it as g(t). Then $g(t) = \check{F}_y(t)e^{-2\pi y \cdot t}$ holds for $y \in B$. Next, we show that $g(t)e^{2\pi y \cdot t} \in L^1(\mathbb{R}^n)$. Let us decompose \mathbb{R}^n into a finite union of non-overlapping polygonal cones, $\Gamma_1, \Gamma_2, \ldots, \Gamma_N$ with their very vertexes at the origin, i.e., $\mathbb{R}^n = \bigcup_{k=1}^N \Gamma_k$. Then $\chi_{\Gamma_k}(t)g(t)e^{2\pi y \cdot t} = \chi_{\Gamma_k}(t)\check{F}_{y_k}(t)e^{-2\pi(y_k-y) \cdot t}$. For any $y_0 \in B$, there exists δ such that $\overline{D(y_0, \delta)} \subset B$. Then for any $y \in D(y_0, \frac{\delta}{4})$ and $y_k \in (y_0 + \Gamma_k)$ satisfying $\frac{3\delta}{4} \leq |y_k - y_0| < \delta$, we get $(y_k - y) \cdot t = (y_k - y_0) \cdot t + (y_0 - y) \cdot t$. Since $y_k - y_0, t \in \Gamma_k$, the angle between the segments $O(y_k - y_0)$ and Ot is less than, say $\frac{\pi}{4}$. Then $(y_k - y) \cdot t \geq \frac{|y_k - y_0|}{\sqrt{2}} |t| - |y_0 - y| |t| \geq (\frac{3}{4\sqrt{2}} - \frac{1}{4})\delta|t| \geq \frac{1}{4}\delta|t|$. Thus, it follows from Hölder's inequality that

$$\int_{\Gamma_k} |g(t)e^{2\pi y \cdot t}| dt \le \int_{\Gamma_k} |\check{F}_{y_k}(t)e^{-\pi \frac{\delta}{4}|t|}| dt \le \|F_{y_k}(x)\|_{L^1(\mathbb{R}^n)} \int_{\Gamma_k} e^{-\pi \frac{\delta}{4}|t|} dt < \infty,$$

which shows that $g(t)e^{2\pi y \cdot t} \in L^1(\Gamma_k)$. Hence $g(t)e^{2\pi y \cdot t} \in L^1(\mathbb{R}^n)$. Together with the relation $g(t) = \check{F}_y(t)e^{-2\pi y \cdot t}$ for $y \in B$, there holds $F(z) = \int_{\mathbb{R}^n} g(t)e^{-2\pi i z \cdot t}$ for all $y \in B$. By letting f(t) = g(-t), we then obtain the desired formula (3) for p = 1 and $z \in T_B$.

Thus, $f(t)e^{-2\pi y \cdot t} \in L^1(\Gamma_k)$ implies that

$$\sup_{t \in \mathbb{R}^{n}} |f(t)|e^{-2\pi y \cdot t} \leq \int_{\mathbb{R}^{n}} |F(x+iy)| dx$$

$$|f(t)|e^{-2\pi y \cdot t}e^{-2\pi \psi(y)} \leq \int_{\mathbb{R}^{n}} |F(x+iy)|e^{-2\pi \psi(y)} dx$$

$$|f(t)|^{s} \int_{B} e^{-2s\pi (y \cdot t+\psi(y))} dy \leq \int_{B} \left(\int_{\mathbb{R}^{n}} |F(x+iy)|e^{-2\pi \psi(y)} dx \right)^{s} dy$$

$$= ||F||_{A^{1,s}(B,\psi)}^{s}, \qquad (8)$$

which implies (4). Next we prove $\operatorname{supp} f \subset U_s(B,\psi)$. Suppose that $t_0 \notin U_s(B,\psi)$, then (1) implies $\int_B e^{-2s\pi(y\cdot t_0+\psi(y))}dy = +\infty$ for $y \in B$. It then follows from (8) that f(t) = 0, which means the support of f, i.e., $\operatorname{supp} f \subset U_s(B,\psi)$.

Next we prove the case $1 . Let <math>B_0 \subseteq B$ be a bounded connected open set, so there exists a positive constant R_0 such that $B_0 \subseteq D(0, R_0)$. Assume that $l_{\varepsilon}(z) = (1 + \varepsilon(z_1^2 + \cdots + z_n^2))^N$, where N is an integer satisfying N > n. Then for $\varepsilon \leq \frac{1}{2R_0^2}$, z = x + iywith $|y| \leq R_0$,

$$|l_{\varepsilon}(z)| = |((1+\varepsilon(|x|^{2}+\cdots+z_{n}^{2}))^{2})^{\frac{N}{2}}|$$

$$= \left(\left(1+\varepsilon(|x|^{2}-|y|^{2})\right)^{2}+4\varepsilon^{2}(x\cdot y)^{2}\right)^{\frac{N}{2}}$$

$$\geq \left(1+\varepsilon(|x|^{2}-|y|^{2})\right)^{N} \ge \left(\frac{1}{2}+\varepsilon|x|^{2}\right)^{N}$$

for $|y| \leq R_0$, i.e., $|l_{\varepsilon}^{-1}(z)| \leq \frac{1}{\left(\frac{1}{2}+\varepsilon|x|^2\right)^N}$. For $F_y(x) = F(x+iy)$, set $F_{\varepsilon,y}(x) = F_y(x)l_{\varepsilon}^{-1}(z)$, then based on Hölder's inequality,

$$\int_{\mathbb{R}^n} |F_{\varepsilon,y}(x)| \, dx \le \left(\int_{\mathbb{R}^n} |F_y(x)|^p \, dx \right)^{\frac{1}{p}} \left(\int_{\mathbb{R}^n} \left| l_{\varepsilon}^{-1}(x+iy) \right|^q \, dx \right)^{\frac{1}{q}} < \infty,$$

where $\frac{1}{p} + \frac{1}{q} = 1$, which implies that $F_{\varepsilon,y}(x) \in L^1(\mathbb{R}^n)$. Then as in the proof for p = 1, $g_{\varepsilon,y}(t) = \check{F}_{\varepsilon,y}(t)e^{-2\pi y \cdot t}$ can be also proved to be independent of $y \in B_0$ when $1 . Put <math>g_{\varepsilon,y}(t) = g_{\varepsilon}(t)$, then $g_{\varepsilon}(t)e^{2\pi y \cdot t} = \check{F}_{\varepsilon,y}(t) \in L^1(\mathbb{R}^n)$.

On the other hand, it is obvious that $F_{\varepsilon,y}(x) \to F_y(x)$ pointwise as $\varepsilon \to 0$. Now we prove that $\check{F}_y(t)e^{-2\pi y \cdot t}$ is also independent of $y \in B_0$. Indeed, for $a, b \in B_0$ and any compact subset $K \subset \mathbb{R}^n$, let $R_1 = \max\{|z| : z \in K\}$,

$$\begin{split} &\left(\int_{K}\left|\check{F}_{a}(t)e^{-2\pi a \cdot t}-\check{F}_{b}(t)e^{-2\pi b \cdot t}\right|^{q}dt\right)^{\frac{1}{q}} \\ &\leq \left(\int_{K}\left|\check{F}_{a}(t)e^{-2\pi a \cdot t}-g_{\varepsilon}(t)\right|^{q}dt\right)^{\frac{1}{q}}+\left(\int_{K}\left|g_{\varepsilon}(t)-\check{F}_{b}(t)e^{-2\pi b \cdot t}\right|^{q}dt\right)^{\frac{1}{q}} \\ &= \left(\int_{K}\left|\check{F}_{a}(t)e^{-2\pi a \cdot t}-\check{F}_{\varepsilon,a}(t)e^{-2\pi a \cdot t}\right|^{q}dt\right)^{\frac{1}{q}}+\left(\int_{K}\left|\check{F}_{\varepsilon,b}(t)e^{-2\pi b \cdot t}-\check{F}_{b}(t)e^{-2\pi b \cdot t}\right|^{q}dt\right)^{\frac{1}{q}} \\ &\leq e^{2\pi R_{0}R_{1}}\left(\left(\int_{K}\left|\check{F}_{a}(t)-\check{F}_{\varepsilon,a}(t)\right|^{q}dt\right)^{\frac{1}{q}}+\left(\int_{K}\left|\check{F}_{\varepsilon,b}(t)-\check{F}_{b}(t)\right|^{q}dt\right)^{\frac{1}{q}}\right) \\ &\leq e^{2\pi R_{0}R_{1}}\left(\left(\int_{\mathbb{R}^{n}}\left|F_{a}(t)-F_{\varepsilon,a}(t)\right|^{p}dt\right)^{\frac{1}{p}}+\left(\int_{\mathbb{R}^{n}}\left|F_{\varepsilon,b}(t)-F_{b}(t)\right|^{p}dt\right)^{\frac{1}{p}}\right) \\ &\rightarrow 0, \end{split}$$

as $\varepsilon \to 0$. Hence we obtain that $\check{F}_a(t)e^{-2\pi a \cdot t} = \check{F}_b(t)e^{-2\pi b \cdot t}$ almost everywhere on \mathbb{R}^n and write it as g(t). Then we have $g(t) = \check{F}_y(t)e^{-2\pi y \cdot t}$.

Next, we show that $g(t)e^{2\pi y \cdot t} \in L^1(\mathbb{R}^n)$. As in the proof for p = 1, let $\mathbb{R}^n = \bigcup_{k=1}^N \Gamma_k$ and $\overline{D(y_0, \delta)} \subset B_0$. Then for any $y \in D(y_0, \frac{\delta}{4})$ and $y_k \in (y_0 + \Gamma_k)$ satisfying $\frac{3\delta}{4} \leq |y_k - y_0| < \delta$, we have

$$(y_k - y) \cdot t \ge \frac{|y_k - y_0|}{\sqrt{2}} |t| - |y_0 - y||t| \ge (\frac{3}{4\sqrt{2}} - \frac{1}{4})\delta|t| \ge \frac{1}{4}\delta|t|$$

for $y_k - y_0, t \in \Gamma_k$. Thus, from Hölder's inequality

$$\int_{\Gamma_k} |g(t)e^{2\pi y \cdot t}|dt \leq \int_{\Gamma_k} |\check{F}_{y_k}(t)e^{-\pi\frac{\delta_0}{4}|t|}|dt \leq \left(\int_{\Gamma_k} |\check{F}_{y_k}(t)|^p dt\right)^{\frac{1}{p}} \left(\int_{\Gamma_k} |e^{-q\pi\frac{\delta_0}{4}|t|}|dt\right)^{\frac{1}{q}} < \infty,$$

which shows that $g(t)e^{2\pi y \cdot t} \in L^1(\Gamma_k)$ and the function G(z) defined by

$$G(z) = \int_{\mathbb{R}^n} g(t) e^{-2\pi i (x+iy) \cdot t} dt$$

is holomorphic in the tube domain $T_{D(y_0,\delta)}$.

Now we can prove that, for $y \in B_0$,

$$\lim_{\varepsilon \to 0} \int_{\mathbb{R}^n} g_{\varepsilon}(t) e^{-2\pi i (x+iy) \cdot t} dt = \int_{\mathbb{R}^n} g(t) e^{-2\pi i (x+iy) \cdot t} dt.$$

In fact, if $y \in B_0$,

$$\begin{aligned} \left| \int_{\mathbb{R}^{n}} (g_{\varepsilon}(t) - g(t)) e^{-2\pi i (x+iy) \cdot t} dt \right| \\ &\leq \int_{\mathbb{R}^{n}} \left| \left(\check{F}_{\varepsilon,y}(t) e^{-2\pi y \cdot t} - \check{F}_{y}(t) e^{-2\pi y \cdot t} \right) e^{2\pi i z \cdot t} \right| dt \\ &= \sum_{k=1}^{n} \int_{\Gamma_{k}} \left| \left(\check{F}_{\varepsilon,y_{k}}(x) - \check{F}_{y_{k}}(x) \right) e^{-2\pi i (y_{k}-y) \cdot t} \right| dt \\ &\leq \sum_{k=1}^{n} \left(\int_{\Gamma_{k}} \left| \check{F}_{\varepsilon,y_{k}}(x) - \check{F}_{y_{k}}(x) \right|^{q} dt \right)^{\frac{1}{q}} \left(\int_{\Gamma_{k}} e^{-p\pi \frac{\delta_{0}}{4} |t|} dt \right)^{\frac{1}{p}} \\ &\leq C_{\delta_{0}} \sum_{k=1}^{n} \left(\int_{\Gamma_{k}} |F_{\varepsilon,y_{k}}(x) - F_{y_{k}}(x)|^{p} dt \right)^{\frac{1}{p}} \\ &\rightarrow 0 \end{aligned}$$

when $\varepsilon \to 0$, where $C_{\delta_0}^p = \int_{\mathbb{R}^n} e^{-p\pi \frac{\delta_0}{4}|t|} dt$. It follows that $\lim_{\varepsilon \to 0} F_\varepsilon(z) = G(z)$. Combining with $\lim_{\varepsilon \to 0} F_\varepsilon(z) = F(z)$, we can state G(z) = F(z) for $y \in B_0$. Then there exists a measurable function g(t) such that $F(z) = \int_{\mathbb{R}^n} g(t)e^{-2\pi i z \cdot t} dt$ holds for $y \in B_0$. Since B is connected, we can choose a sequence of bounded connected open set $\{B_k\}$ such that $B_0 \subset B_1 \subset \cdots$ and $B = \bigcup_{k=0}^{\infty} B_k$. Together with the fact that $g(t) = \check{F}_y(t)e^{-2\pi y \cdot t}$ is independent of $y \in B_k$, then $\check{F}_{y_l}(t)e^{-2\pi y_l \cdot t} = \check{F}_{y_j}(t)e^{-2\pi y \cdot t} = \check{F}_y(t)e^{-2\pi y \cdot t}$ for $l \neq j$, $y_l \in B_l$, $y_j \in B_j$ and $y \in B_0$. Hence $g(t)e^{2\pi y \cdot t} = \check{F}_y(t)$ holds for $y \in B_k$, $k = 0, 1, 2, \ldots$. In other words, $f(z) = \int_{\mathbb{R}^n} g(t)e^{-2\pi i z \cdot t} dt$ holds for all $y \in B$. By letting f(t) = g(-t), we obtain the desired representation $F(z) = \int_{\mathbb{R}^n} f(t)e^{2\pi i z \cdot t} dt$ for $y \in B$ when 1 .

For $\frac{1}{p} + \frac{1}{q} = 1$, based on the Hausdorff-Young Inequality,

$$\left(\int_{\mathbb{R}^n} |f(t)e^{-2\pi y \cdot t}|^q dt\right)^{\frac{1}{q}} \le \left(\int_{\mathbb{R}^n} |F(x+iy)|^p dx\right)^{\frac{1}{p}},\tag{9}$$

then

$$\left(\left(\int_{\mathbb{R}^n} |f(t)e^{-2\pi y \cdot t}|^q dt\right)^{\frac{p}{q}} e^{-2p\pi\psi(y)} dy\right)^s \le \left(\int_{\mathbb{R}^n} |F(x+iy)e^{-2\pi\psi(y)}|^p dx\right)^s.$$

Performing integral about $y \in B$ on both sides, we get

$$\int_B \left(\left(\int_{\mathbb{R}^n} |f(t)e^{-2\pi y \cdot t}|^q dt \right)^{\frac{p}{q}} e^{-2p\pi\psi(y)} \right)^s dy \le \int_B \left(\int_{\mathbb{R}^n} |F(x+iy)e^{-2\pi\psi(y)}|^p dx \right)^s dy$$

and

$$\int_{B} \left(\left(\int_{\mathbb{R}^{n}} |f(t)e^{-2\pi y \cdot t}|^{q} dt \right)^{\frac{p}{q}} e^{-2p\pi\psi(y)} \right)^{s} dy \leq \|F\|_{A^{p,s}(B,\psi)}^{sp}.$$
(10)

As a result, formulas (3) and (5) hold for $1 . Now we prove that <math>\operatorname{supp} f \subset U_{sp}(B, \psi)$ when $0 < s(p-1) \le 1$. For $0 < s(p-1) \le 1$, we have $\frac{q}{sp} \ge 1$. Then Minkowski's inequality and (10) imply that

$$\int_{\mathbb{R}^n} |f(t)|^q \left(\int_B e^{-2\pi p s(y \cdot t + \psi(y))} dy \right)^{\frac{q}{ps}} dt \le \|F\|_{A^{p,s}(B,\psi)}^q < \infty.$$
(11)

Consequently, It follows from (11) and (1) that f(t) = 0 for almost every $t \notin U_{sp}(B, \psi)$. Therefore, $\operatorname{supp} f \subset U_{sp}(B, \psi)$.

In order to prove Theorem 2, we first introduce a lemma.

LEMMA 1. Suppose that $F(z) \in A^{p,s}(B,\psi)$, where $0 and <math>0 < s \le \infty$, then for $y_0 \in B$ and positive constant δ such that $D_n(y_0,\delta) \subset B$, there exist constants N > 1 and $C_{n,N,p,s}$ depending on n, N, p, s such that

$$|F(z)| \le C_{n,N,p,s} \delta^{-\frac{n}{p}(1+\frac{1}{s})} e^{2\pi\psi_{\delta}(y_0)},$$
(12)

where $\psi_{\delta}(y_0) = \max\{\psi(\eta) : |\eta - y_0| \le \delta\}.$

Proof. For $y_0 \in B$, there exists $\delta > 0$ such that $B_{\delta} = D(y_0, \delta) \subset B$. Then for $F(z) = F(x+iy) \in A^{p,s}(B,\psi)$, based on the subharmonic properties of $|F(z)|^t$, we have

$$|F(z)|^t \leq \frac{1}{\Omega_{2n}\delta^{2n}} \int_{D_{2n}(z,\delta)} |F(\xi+i\eta)|^t d\xi d\eta \leq \frac{1}{\Omega_{2n}\delta^{2n}} \int_{D_n(y_0,\delta)} \left(\int_{D_n(x,\delta)} |F(\xi+i\eta)|^t d\xi \right) d\eta$$

for $y \in B_{\delta}$, where Ω_k is the volume of k-dimensional unit ball $D_k(0, 1)$ centered at the origin with radius 1, k = n, 2n. Let $p_1 = N = \frac{p}{t} > \max\{1, \frac{1}{s}\}$ and $\frac{1}{p_1} + \frac{1}{q_1} = 1$. Hölder's Inequality implies that

$$|F(z)|^{t} \leq \frac{1}{\Omega_{2n}\delta^{2n}} \int_{D_{n}(y_{0},\delta)} \left(\int_{D_{n}(x,\delta)} |F(\xi+i\eta)|^{p} d\xi \right)^{\frac{1}{p_{1}}} d\eta \left(\int_{D_{n}(x,\delta)} 1^{q_{1}} d\xi \right)^{\frac{1}{q_{1}}} \\ = \frac{(\delta^{n}\Omega_{n})^{\frac{1}{q_{1}}}}{\delta^{2n}\Omega_{2n}} \int_{D_{n}(y_{0},\delta)} \left(\int_{D_{n}(x,\delta)} |F(\xi+i\eta)|^{p} d\xi \right)^{\frac{1}{p_{1}}} d\eta.$$

For $0 < s < \infty$, let $p_2 = sN$. Then $p_2 > 1$. Again, by Hölder's Inequality, for $\frac{1}{p_2} + \frac{1}{q_2} = 1$,

$$\begin{split} |F(z)|^{t} &\leq \frac{(\delta^{n}\Omega_{n})^{\frac{1}{q_{1}}}}{\delta^{2n}\Omega_{2n}} \left(\int_{D_{n}(y_{0},\delta)} \left(\int_{D_{n}(x,\delta)} |F(\xi+i\eta)|^{p} d\xi \right)^{s} d\eta \right)^{\frac{1}{p_{2}}} \left(\int_{D_{n}(y_{0},\delta)} 1^{q_{2}} d\eta \right)^{\frac{1}{q_{2}}} \\ &\leq \frac{(\delta^{n}\Omega_{n})^{\frac{1}{q_{1}}+\frac{1}{q_{2}}}}{\delta^{2n}\Omega_{2n}} \left(\int_{D_{n}(y_{0},\delta)} \left(\int_{D_{n}(x,\delta)} |F(\xi+i\eta)e^{-2\pi\psi(\eta)}|^{p} d\xi \right)^{s} e^{2sp\pi\psi(\eta)} d\eta \right)^{\frac{1}{p_{2}}} \\ &\leq \frac{(\delta^{n}\Omega_{n})^{2-\frac{1}{N}(1+\frac{1}{s})}e^{2\frac{sp}{p_{2}}\pi\psi_{\delta}(y_{0})}}{\delta^{2n}\Omega_{2n}} \left(\int_{D_{n}(y_{0},\delta)} \left(\int_{D_{n}(y_{0},\delta)} |F(\xi+i\eta)e^{-2\pi\psi(\eta)}|^{p} d\xi \right)^{s} d\eta \right)^{\frac{1}{p_{2}}} \\ &\leq \frac{(\delta^{n}\Omega_{n})^{2-\frac{1}{N}(1+\frac{1}{s})}e^{2\frac{sp}{p_{2}}\pi\psi_{\delta}(y_{0})}}{\delta^{2n}\Omega_{2n}} \left(\int_{B} \left(\int_{\mathbb{R}^{n}} |F(\xi+i\eta)e^{-2\pi\psi(\eta)}|^{p} d\xi \right)^{s} d\eta \right)^{\frac{1}{p_{2}}}, \end{split}$$

where $\psi_{\delta}(y_0) = \max\{\psi(\eta) : |\eta - y_0| \le \delta\}$. Hence,

$$\begin{aligned} |F(z)| &\leq \left(\frac{\delta^{-\frac{n}{N}(1+\frac{1}{s})}\Omega_{n}^{2-\frac{1}{N}(1+\frac{1}{s})}e^{2\frac{sp}{p_{2}}\pi\psi_{\delta}(y_{0})}}{\Omega_{2n}} \right)^{\frac{1}{t}} \left(\int_{B} \left(\int_{\mathbb{R}^{n}} |F(\xi+i\eta)e^{-2\pi\psi(\eta)}|^{p}d\xi \right)^{s}d\eta \right)^{\frac{1}{tp_{2}}} \\ &\leq \frac{\Omega_{n}^{\frac{2N}{p}-\frac{1}{p}(1+\frac{1}{s})}}{\Omega_{2n}^{\frac{N}{p}}\delta^{\frac{n}{p}(1+\frac{1}{s})}}e^{2\frac{sp}{tp_{2}}\pi\psi_{\delta}(y_{0})} \left(\int_{B} \left(\int_{\mathbb{R}^{n}} |F(\xi+i\eta)e^{-2\pi\psi(\eta)}|^{p}d\xi \right)^{s}d\eta \right)^{\frac{1}{sp}\frac{sp}{tp_{2}}}. \end{aligned}$$

Since $\frac{sp}{tp_2} = 1$, by letting $C_{n,N,p,s} = \frac{\Omega_n^{\frac{2N}{p} - \frac{1}{p}(1 + \frac{1}{s})}}{\Omega_{2n}^{\frac{N}{p}}} \|F(z)\|_{A^{p,s}(B,\psi)}$, we obtain the desired inequality

$$|F(z)| \le C_{n,N,p,s} \delta^{-\frac{n}{p}(1+\frac{1}{s})} e^{2\pi\psi_{\delta}(y_0)}.$$

While $s = \infty$, for $p_2 = sN = \infty$, we have

$$|F(z)|^t \le \frac{(\delta^n \Omega_n)^{2-\frac{1}{N}}}{\delta^{2n} \Omega_{2n}} \sup_{\eta \in D_n(y,\delta)} \left| \int_{D_n(x,\delta)} |F(\xi + i\eta)|^p d\xi \right|^{\frac{1}{p}}.$$

Then

$$|F(z)| \leq \frac{(\delta^{n}\Omega_{n})^{(2-\frac{1}{N})\frac{N}{p}}}{(\delta^{2n}\Omega_{2n})^{\frac{N}{p}}} e^{2\pi\psi_{\delta}(y_{0})} \sup_{\eta\in D_{n}(y,\delta)} \left| \left(\int_{D_{n}(x,\delta)} |F(\xi+i\eta)|^{p} d\xi \right)^{\frac{1}{p}} e^{-2\pi\psi(y)} \right|$$
$$= \frac{\Omega_{n}^{\frac{2N}{p}-\frac{1}{p}}}{\Omega_{2n}^{\frac{N}{p}}} \delta^{-\frac{n}{p}} e^{2\pi\psi_{\delta}(y_{0})} ||F(z)||_{A^{p,\infty}(B,\psi)}.$$

Obviously, the inequality (12) is also applicable in the case $s = \infty$.

Now we are ready to prove Theorem 2.

Proof of Theorem 2. For $y_0 \in B$, there exists $\delta > 0$ such that $B_{\delta} = D(y_0, \delta) \subset B$. Then for $F(z) \in A^{p,s}(B, \psi)$ and any $y \in B_{\delta}$, it follows from Lemma 1 that

$$|F(z)| \le C_{n,N,p,s} \delta^{-\frac{n}{p}(1+\frac{1}{s})} e^{2\pi\psi_{\delta}(y_0)}.$$

Thus,

$$\int_{\mathbb{R}^n} |F(z)|^2 dx = \int_{\mathbb{R}^n} |F(z)|^{p+2-p} dx \le C_{n,N,p,s}^{2-p} \delta^{-\frac{n(2-p)}{p}(1+\frac{1}{s})} e^{2(2-p)\pi\psi_{\delta}(y_0)} \int_{\mathbb{R}^n} |F(z)|^p dx.$$

Therefore,

$$\begin{split} & \int_{\mathbb{R}^n} |F(z)|^2 e^{-4\pi\psi_{\delta}(y_0)} dx \\ & \leq \quad C_{n,N,p,s}^{2-p} \delta^{-\frac{n(2-p)}{p}(1+\frac{1}{s})} e^{2(2-p)\pi\psi_{\delta}(y_0)} \int_{\mathbb{R}^n} |F(z)e^{-2\pi\psi(y)}|^p dx e^{2p\pi\psi(y)}e^{-4\pi\psi_{\delta}(y_0)} \\ & \leq \quad C_{n,N,p,s}^{2-p} \delta^{-\frac{n(2-p)}{p}(1+\frac{1}{s})} e^{2(2-p)\pi\psi_{\delta}(y_0)} \int_{\mathbb{R}^n} |F(z)e^{-2\pi\psi(y)}|^p dx e^{2(p-2)\pi\psi_{\delta}(y_0)} \\ & = \quad C_{n,N,p,s}^{2-p} \delta^{-\frac{n(2-p)}{p}(1+\frac{1}{s})} \int_{\mathbb{R}^n} |F(z)e^{-2\pi\psi(y)}|^p dx. \end{split}$$

Taking integral with respect to y to both sides of the inequality, we have

$$\int_{B_{\delta}} \left(\int_{\mathbb{R}^{n}} |F(z)|^{2} e^{-4\pi\psi_{\delta}(y_{0})} dx \right)^{s} dy \leq C_{n,N,p,s}^{(2-p)s} \delta^{-\frac{n(2-p)(1+s)}{p}} \int_{B_{\delta}} \left(\int_{\mathbb{R}^{n}} |F(z)e^{-2\pi\psi(y)}|^{p} dx \right)^{s} dy,$$

which concludes that $F \in A^{2,s}(B_{\delta}, \psi_{\delta})$. Similarly, we can prove that

$$\int_{B_{\delta}} \left(\int_{\mathbb{R}^{n}} |F| e^{-2\pi\psi_{\delta}(y_{0})} dx \right)^{s} dy \leq C_{n,N,p,s}^{(1-p)s} \delta^{-\frac{n(1-p)(1+s)}{p}} \|F(z)\|_{A^{1,s}(B_{\delta},\psi)}^{sp}.$$
 (13)

Then $F(z) \in A^{1,s}(B_{\delta}, \psi_{\delta}).$

Following the proof of the case p = 1 in Theorem 1, there exists a continuous function f(t) such that $F_y(x) = \int_{\mathbb{R}^n} f(t)e^{2\pi i z \cdot t} dt$ holds for $y \in B_\delta$ and $f(t) = \hat{F}_y(t)e^{2\pi y \cdot t}$ is independent of $y \in B$. Together with the fact that $f(t)e^{-2\pi y \cdot t} \in L^1(\mathbb{R}^n)$ for all $y \in B$, we see that (3) holds for all $y \in B$. This completes the proof of Theorem 2.

Before the proof of Corollary 1, we introduce the following lemma.

LEMMA 2. Assume that Γ is a regular open convex cone of \mathbb{R}^n . Let $\psi \in C(\Gamma)$ satisfy (6), then $U_{\alpha}(\psi,\Gamma) \subset \Gamma^* + \overline{D(0,R_{\psi})}$, where $U_{\alpha}(\psi,\Gamma)$ is defined by (1) for $0 < \alpha < \infty$ and by (2) for $\alpha = \infty$.

Proof. For $t_0 \notin \Gamma^* + \overline{D(0, R_{\psi})}$, there exist $\varepsilon > 0$ and $\xi \in \Gamma^*$ such that $d(t_0, \Gamma^*) = |\xi - t_0| \ge R_{\psi} + 3\varepsilon$ and $\xi \cdot (t_0 - \xi) = 0$. Then for any $\tilde{t} \in \Gamma^*$,

$$(\tilde{t} - t_0) \cdot \frac{(\xi - t_0)}{|\xi - t_0|} \ge |\xi - t_0|.$$

Hence $\tilde{t} \cdot (\xi - t_0) = (\tilde{t} - t_0 + t_0 - \xi + \xi) \cdot (\xi - t_0) \ge |\xi - t_0|^2 - |\xi - t_0|^2 = 0$, which means $\xi - t_0 \in \overline{\Gamma}$. For any $\delta > 0$, it follows from (6) that there exists ρ_0 such that $\psi(y) \le (R_{\psi} + \delta)|y|$ for $|y| \ge \rho_0$. Let

 $e_{0} = \frac{\xi - t_{0}}{|\xi - t_{0}|} \in \overline{\Gamma} \cap \partial D(0, 1), \text{ then for any } \varepsilon_{1} > 0, \text{ we can find an } e_{1} \in \Gamma \text{ with } |e_{1}| = 1 \text{ such that } |e_{1} - e_{0}| < \varepsilon_{1}, \text{ which means there exists a positive constant } \delta_{1} < \varepsilon_{1} \text{ such that } D(e_{1}, \delta_{1}) \subset \Gamma.$ Thus, for any $e \in D(e_{1}, \delta_{1})$ with $|e_{1}| = 1$, we have $|e - e_{0}| \leq |e - e_{1}| + |e_{1} - e_{0}| < 2\varepsilon_{1}.$ Choose ε_{1} satisfying $2\varepsilon_{1}|t_{0}| \leq \varepsilon$ and let $\Gamma_{1} = \{y = \rho e : \rho > 0 \text{ and } e \in D(e_{1}, \delta) \cap \partial D(0, 1)\} \subset \Gamma.$ Then for any $y \in \Gamma_{1}, -\rho e \cdot t_{0} = \rho(-e + e_{0} - e_{0}) \cdot t_{0} \geq \rho(-2\varepsilon_{1}|t_{0}| + |\xi - t_{0}|) \geq \rho(R_{\psi} + 2\varepsilon)$ and

$$\int_{\Gamma} e^{-2\pi\alpha(t_0 \cdot y + \psi(y))} dy \geq \int_{\Gamma \cap \{|y| \ge \rho_0\}} e^{-2\pi\alpha(t_0 \cdot y + (R_\psi + \delta)|y|)} dy$$

$$\geq \int_{\rho_0}^{\infty} \rho^{n-1} d\rho \int_{\partial D(0,1) \cap D(e_1,\delta_1)} e^{2\pi\alpha\rho(2\varepsilon - \delta)} d\sigma(\zeta) = +\infty,$$

which implies $t_0 \notin U_{\alpha}(\psi, \Gamma)$. Therefore, $U_{\alpha}(\psi, \Gamma) \subset \Gamma^* + \overline{D(0, R_{\psi})}$.

Now we prove Corollary 1.

Proof of Corollary 1. For $y_0 \in \Gamma$, there exists δ such that $D(y_0, \delta) \subset \Gamma$. It follows from Theorem 2 that there exists f(t) such that (3) holds for $y \in D(y_0, \delta)$. Since Γ is connected, (3) also holds for all $y \in \Gamma$. Applying the methods in the proof of Theorem 1 for p = 1, we obtain that such an f(t) is supported in $U_s(\Gamma, \psi_{\delta})$. Combing with Lemma 2, we have $\operatorname{supp} f \subset U_s(\Gamma, \psi_{\delta}) \subset \Gamma^* + \overline{D(0, R_{\psi_{\delta}})}$, where

$$R_{\psi_{\delta}} = \lim_{y \in \Gamma, y \to \infty} \frac{\psi_{\delta}(y)}{|y|}.$$

Since $R_{\psi_{\delta}} = R_{\psi}$ for any $y \in \Gamma$, we see that $U_s(\Gamma, \psi_{\delta})$ is also a subset of $\Gamma^* + \overline{D(0, R_{\psi})}$. Hence, supp $f \subset \Gamma^* + \overline{D(0, R_{\psi})}$.

Now we show that $|f(t)| \left(\int_{\Gamma} e^{-2s\pi(y\cdot t+R_{\psi}|y|)} dy \right)^{\frac{1}{s}}$ is slowly increasing. For $y_0, y \in \Gamma$, $y_0 + y \in \Gamma$, $F_{y_0}(z) = F(x + i(y + y_0)) \in A^{p,s}(\Gamma, \psi)$. As in Theorem 1, we have $f(t) = g(-t) = \check{F}_{y_0+y}(-t)e^{2\pi(y_0+y)\cdot t}$. Due to the relation $R_{\psi} = \lim_{y \in B, y \to \infty} \frac{\psi(y)}{|y|}$, we have $\psi_{\delta}(y) \leq R_{\psi}(1+|y_0|+|y|)$, where R_{ψ} is a positive constant independent of $y_0, y \in \Gamma$. Then

$$\begin{aligned} |f(t)| &= |\check{F}_{y_0+y}(-t)e^{2\pi(y_0+y)\cdot t}| = \left| \int_{\mathbb{R}^n} F_{y_0+y}(x)e^{-2\pi ix\cdot t}e^{-2\pi\psi_{\delta}(y)}dx \right| e^{2\pi(\psi_{\delta}(y)+(y_0+y)\cdot t)} \\ &\leq \int_{\mathbb{R}^n} |F_{y_0}(z)|e^{-2\pi\psi_{\delta}(y)}dx e^{2\pi(R_{\psi}(1+|y_0|+|y|)+(y_0+y)\cdot t)}. \end{aligned}$$

Combining with (13), it follows that

$$\left(\int_{\Gamma} |f(t)|^{s} e^{-2s\pi(y\cdot t + R_{\psi}|y|)} dy \right)^{\frac{1}{s}} \leq \left(\int_{\Gamma} \left(\int_{\mathbb{R}^{n}} |F_{y_{0}}(z)| e^{-2\pi\psi_{\delta}(y)} dx \right)^{s} dy \right)^{\frac{1}{s}} e^{2\pi(R_{\psi}(1+|y_{0}|)+y_{0}\cdot t)}$$

$$\leq C_{n,N,p,s}^{1-p} \delta^{-\frac{n(1-p)(1+s)}{sp}} \|F_{y_{0}}\|_{A^{1,s}(B,\psi)}^{p} e^{2\pi(R_{\psi}(1+|y_{0}|)+y_{0}\cdot t)}$$

$$= C \exp\{J(y_{0},t)\},$$

where $C = C_{n,N,p,s}^{1-p} \|F_{y_0}\|_{A^{1,s}(\Gamma,\psi)}^p$ and $J(y_0) = -\frac{n(1-p)(1+s)}{sp} \log \delta + 2\pi (R_{\psi}(1+|y_0|)+y_0 \cdot t)$. Let $J(t) = \inf\{J(y_0,t) : y_0 \in \Gamma\}$, then

$$|f(t)| \left(\int_{\Gamma} e^{-2s\pi(y\cdot t + R_{\psi}|y|)} dy \right)^{\frac{1}{s}} \le C \exp\{J(t)\}$$

Take $y_0 = \rho v$ with $\rho > 0$ and a fixed $v \in \Gamma$ with |v| = 1, then $\delta = d(\rho v, \partial \Gamma)/2 = \rho \varepsilon$, where $\varepsilon = d(v, \partial \Gamma)/2$. Therefore,

$$J(t) = \inf_{\rho > 0} \left\{ -\frac{n(1-p)(1+s)}{sp} \log(\varepsilon\rho) + 2\pi R_{\psi}(1+\rho) + 2\pi\rho|t| \right\},\$$

in which the infimum can be attained when $\rho = \frac{n(1-p)(1+s)}{2sp\pi(R_{\psi}+|t|)}$. It follows that

$$J(t) \le 2\pi R_{\psi} + n\left(\frac{1}{p} - 1\right)\left(\frac{1}{s} + 1\right)\left(1 - \log\varepsilon - \log n\left(\frac{1}{p} - 1\right)\left(\frac{1}{s} + 1\right) + \log 2\pi (R_{\psi} + |t|)\right).$$

Thus, there exists a positive constant $M_{n,p,s,v}$ such that

$$|f(t)| \left(\int_{\Gamma} e^{-2s\pi(y\cdot t + R_{\psi}|y|)} dy \right)^{\frac{1}{s}} \le C e^{J(t)} \le M_{n,p,s,v} (1 + |t|)^{n(\frac{1}{p} - 1)(\frac{1}{s} + 1)}.$$

The proof is complete.

Proof of Theorem 3. We first prove the case when $2 . Since <math>\Gamma$ is a regular open convex cone, $\operatorname{int}\Gamma \neq \emptyset$, where $\operatorname{int}\Gamma$ is denoted as the interior of Γ . Then for $y \in \Gamma$, we can find a basis $\{e_j\} \subset \operatorname{int}\Gamma^*$ such that $y = \sum_{j=1}^n e_j y_j$ and $e_j \cdot y \ge 0$. For $\varepsilon > 0$, let $l_{\varepsilon}(z) = \left(\prod_{j=1}^n (1 - i\varepsilon e_j \cdot z)\right)^{2N}$ with $N > \frac{n}{2} \left(1 - \frac{1}{p}\right)$ and choose two positive constant A, Bsuch that $B|x|^2 \le \varepsilon^2 \sum_{j=1}^n (e_j \cdot x)^2 \le A|x|^2$ for all $x \in \mathbb{R}^n$. Thus,

$$\begin{aligned} |l_{\varepsilon}(z)| &= \left(\prod_{j=1}^{n} |1 - i\varepsilon e_j \cdot z|^2\right)^N = \left(\prod_{j=1}^{n} \left((1 + \varepsilon e_j \cdot y)^2 + \varepsilon^2 (e_1 \cdot x)^2\right)\right)^N \\ &\geq \left(\prod_{j=1}^{n} \left(1 + \varepsilon^2 (e_j \cdot x)^2\right)\right)^N \ge \left(1 + \varepsilon^2 \sum_{j=1}^{n} (e_j \cdot x)^2\right)^N \ge \left(1 + \varepsilon^2 B |x|^2\right)^N, \end{aligned}$$

i.e., $|l_{\varepsilon}^{-1}(z)| \leq (1 + \varepsilon^2 B|x|^2)^{-N}$. For $F(x + iy) \in A^{p,s}(\Gamma, \psi)$, $F_y(x) = F(x + iy) \in L^p(\mathbb{R}^n)$ as a function of x. Let $F_{\varepsilon}(z) = F_{\varepsilon,y}(x) = F_y(x)l_{\varepsilon}^{-1}(z)$, then $F_{\varepsilon,y}(x) \in L^1(\mathbb{R}^n) \cap L^2(\mathbb{R}^n)$. Indeed, Hölder's inequality implies that

$$\int_{\mathbb{R}^n} |F_{\varepsilon,y}(x)| dx \le \left(\int_{\mathbb{R}^n} |F_y(x)|^p dx \right)^{\frac{1}{p}} \left(\int_{\mathbb{R}^n} |l_{\varepsilon}^{-1}(x+iy)|^q dx \right)^{\frac{1}{q}} \le C_{1,\varepsilon} \|F_y\|_{L^p(\mathbb{R}^n)}$$
(14)

and

$$\int_{\mathbb{R}^n} |F_{\varepsilon,y}(x)|^2 dx \le \left(\int_{\mathbb{R}^n} |F_y(x)|^{\frac{p}{2}} dx \right)^{\frac{2}{p}} \left(\int_{\mathbb{R}^n} |l_{\varepsilon}^{-1}(x+iy)|^{\frac{p}{p-2}} dx \right)^{\frac{p-2}{p}} \le C_{2,\varepsilon} \|F_y\|_{L^p(\mathbb{R}^n)},$$

where $C_{1,\varepsilon} = \left(\int_{\mathbb{R}^n} \frac{dx}{(1+\varepsilon^2 B|x|^2)^{qN}}\right)^{\frac{1}{q}} < \infty, \ C_{2,\varepsilon} = \left(\int_{\mathbb{R}^n} \frac{dx}{(1+\varepsilon^2 B|x|^2)^{\frac{p}{p-2}N}}\right)^{\frac{p-p}{p}} < \infty.$

As the proof of p = 1 in Theorem 1, we can show $g_{\varepsilon}(t)e^{2\pi y \cdot t} = \check{F}_{\varepsilon,y}(t) \in L^1(\mathbb{R}^n)$. Thus,

$$g_{\varepsilon}(t)e^{2\pi y \cdot t} = \int_{\mathbb{R}^n} F_{\varepsilon,y}(x)e^{2\pi i x \cdot t} dx, \qquad (15)$$

then $|g_{\varepsilon}(t)|e^{2\pi y \cdot t} \leq \int_{\mathbb{R}^n} |F_{\varepsilon,y}(x)| dx$. Together with (14), there hold

$$\begin{aligned} |g_{\varepsilon}(t)|e^{2\pi(y\cdot t-\psi(y))} &\leq C_{1,\varepsilon} \left(\int_{\mathbb{R}^n} |F(x+iy)e^{-2\pi\psi(y)}|^p dx \right)^{\frac{1}{p}}, \\ \int_{\Gamma} |g_{\varepsilon}(t)|^{sp} e^{2sp\pi(y\cdot t-\psi(y))} dy &\leq C_{1,\varepsilon} \int_{\Gamma} \left(\int_{\mathbb{R}^n} |F(x+iy)e^{-2\pi\psi(y)}|^p dx \right)^s dy, \\ |g_{\varepsilon}(t)|^{sp} \int_{\Gamma} e^{2sp\pi(y\cdot t-\psi(y))} dy &\leq C_{1,\varepsilon} ||F||_{A^{p,s}(\Gamma,\psi)}^{sp}. \end{aligned}$$

Now we prove that $\operatorname{supp} g_{\varepsilon}(t) \subset -U_{ps}(\Gamma, \psi)$. Note that $g_{\varepsilon}(t)$ is continuous in \mathbb{R}^n . Then for $t_0 \notin -U_{ps}(\Gamma, \psi)$, formula (1) shows that $\int_{\Gamma} e^{2ps\pi(y\cdot t_0-\psi(y))}dy = \infty$ for $y \in \Gamma$. It follows from the above inequality that $g_{\varepsilon}(t_0) = 0$ for $t_0 \notin -U_{ps}(\Gamma, \psi)$. As a result, $\operatorname{supp} g_{\varepsilon}(t) \subset -U_{ps}(\Gamma, \psi)$.

Since $g_{\varepsilon}(t)e^{2\pi y \cdot t} \in L^1(\mathbb{R}^n)$, we can rewrite (15) as

$$F_{\varepsilon,y}(x) = \int_{\mathbb{R}^n} g_{\varepsilon}(t) e^{-2\pi i z \cdot t} \chi_{-U_{ps}(\Gamma,\psi)}(t) dt.$$
(16)

Plancherel's Theorem implies that $\int_{\mathbb{R}^n} |g_{\varepsilon}(t)e^{2\pi y \cdot t}|^2 dt = \int_{\mathbb{R}^n} |F_{\varepsilon,y}(x)|^2 dx$. Then based on Fatou's lemma,

$$\int_{\mathbb{R}^n} |g_{\varepsilon}(t)|^2 \le \lim_{y \in \Gamma, y \to 0} \int_{\mathbb{R}^n} |F(x+iy)|^2 dx < \infty.$$

Thus, there exist $g(t) \in L^2(\mathbb{R}^n)$ and a sequence $\{\varepsilon_k\}$ tending to zero as $k \to \infty$ such that $\lim_{k\to\infty} \int_{\mathbb{R}^n} g_{\varepsilon_k}(t)h(t)dt = \int_{\mathbb{R}^n} g(t)h(t)dt \text{ for } h(t) \in L^2.$ In fact, for $t \in -U_{ps}(\Gamma, \psi)$, lemma 2 implies that $t \in -\Gamma_k^* + \overline{D(0, R_{\psi})}$. Then t can always be written as $t_1 + t_2$ with $t_1 \in -\Gamma_k^*$ and $|t_2| < R_{\psi}$. Hence, for $y \in \Gamma$,

$$y \cdot t = y \cdot (t_1 + t_2) \le -|t_1|k + |t_2||y| \le -(|t| - |t_2|)k + R_{\psi}|t| \le (R_{\psi} - k)|t| + R_{\psi}k,$$

implying that $\int_{\mathbb{R}^n} |e^{2\pi y \cdot t} \chi_{-U_{ps}(B_k,\psi)}(t)|^2 dt < \infty$. Therefore, on the right hand side of (16),

$$\lim_{k \to \infty} \int_{\mathbb{R}^n} g_{\varepsilon_k}(t) e^{-2\pi i z \cdot t} \chi_{-U_{ps}(\Gamma,\psi)}(t) dt = \int_{\mathbb{R}^n} g(t) e^{-2\pi i z \cdot t} \chi_{-U_{ps}(\Gamma,\psi)}(t) dt$$

for $e^{2\pi y \cdot t} \chi_{-U_{ps}(\Gamma,\psi)}(t) \in L^2(\mathbb{R}^n)$. Whilst it is obvious that $F_{\varepsilon}(z) \to F(z)$ when $\varepsilon \to 0$. Sending k to ∞ on both sides of (16) and letting f(t) = g(-t), we obtain that $f \in L^2(\mathbb{R}^n)$ and the support supp f is contained in $U_{ps}(\Gamma,\psi)$, as well as the desired representation (3) holds for all $z \in T_{\Gamma}$.

We now prove the case when $p = \infty$. For $z = (z_1, \ldots, z_n) \in T_{\Gamma}$ and $\varepsilon > 0$, we can also construct a function $F_{\varepsilon,y}(x) = F_{\varepsilon}(z) = F_y(x)l_{\varepsilon}^{-1}(z)$, where $l_{\varepsilon}(z) = \left(\prod_{j=1}^n (1 - i\varepsilon e_j \cdot z)\right)^{2N}$ with $N > \frac{n}{2}$. Then

$$\int_{\mathbb{R}^n} |F_{\varepsilon,y}(x)| dx \le \sup_{x \in \mathbb{R}^n} |F_y(x)| \int_{\mathbb{R}^n} |l_{\varepsilon}^{-1}(x+iy)| dx \le \widetilde{C}_{1,\varepsilon} \|F_y\|_{L^{\infty}(\mathbb{R}^n)} < \infty$$
(17)

and

$$\int_{\mathbb{R}^n} |F_{\varepsilon,y}(x)|^2 dx \le \sup_{x \in \mathbb{R}^n} |F_y(x)| \int_{\mathbb{R}^n} |l_{\varepsilon}^{-1}(x+iy)|^2 dx \le \widetilde{C}_{2,\varepsilon} ||F_y||_{L^{\infty}}(\mathbb{R}^n) < \infty,$$

where $\widetilde{C}_{1,\varepsilon} = \int_{\mathbb{R}^n} \frac{dx}{(1+\varepsilon^2 B|x|^2)^N}$ and $\widetilde{C}_{2,\varepsilon} = \left(\int_{\mathbb{R}^n} \frac{dx}{(1+\varepsilon^2 B|x|^2)^{2N}}\right)^{\frac{1}{2}} < \infty$. Hence $F_{\varepsilon,y} \in L^1(\mathbb{R}^n) \cap L^2(\mathbb{R}^n)$. In this case, we also have $g_{\varepsilon}(t)e^{2\pi y \cdot t} = \check{F}_{\varepsilon,y}(t) \in L^1(\mathbb{R}^n)$. Then $g_{\varepsilon}(t)e^{2\pi y \cdot t} = \int_{\mathbb{R}^n} F_{\varepsilon,y}(x)e^{2\pi ix \cdot t}dx$. Therefore, together with (17),

$$\begin{aligned} |g_{\varepsilon}(t)|e^{2\pi(y\cdot t-\psi(y))} &\leq \widetilde{C}_{1,\varepsilon} \sup_{x\in\mathbb{R}^n} |F_y(x)|e^{-2\pi\psi(y)},\\ \sup_{y\in\Gamma} |g_{\varepsilon}(t)|e^{2\pi(y\cdot t-\psi(y))} &\leq \widetilde{C}_{1,\varepsilon} \sup_{x\in\mathbb{R}^n,y\in\Gamma} |F(x+iy)|e^{-2\pi\psi(y)}\\ &= \widetilde{C}_{1,\varepsilon} ||F||_{A^{\infty,\infty}(\Gamma,\psi)} < \infty. \end{aligned}$$

Then we can similarly show that $\operatorname{supp} g_{\varepsilon}(t) \subset -U_{\infty}(\Gamma, \psi) \subset -\Gamma^* + \overline{D(0, R_{\psi})}$. Applying the same method for $2 , we obtain the desired formula (3) holds for all <math>z \in T_{\Gamma}$ and the support $\operatorname{supp} f$ is contained in $U_{\infty}(\Gamma, \psi) \subset \Gamma^* + \overline{D(0, R_{\psi})}$.

4 Applications

In [10], denoting by $A^2_{\alpha}(\mathbb{C}^+)$ a weighted Bergman space of functions holomorphic in \mathbb{C}^+ satisfying $||F||^2_{A^2_{\alpha}(\mathbb{C}^+)} = \int_{\mathbb{C}^+} |F(x+iy)|^2 y^{\alpha} dx dy < \infty$, and by $L^2_{\beta}(\mathbb{R}^+)$ the space of complex– valued measurable functions f on \mathbb{R}^+ satisfying $||f||^2_{L^2_{\beta}(\mathbb{R}^+)} = \frac{\Gamma(\beta)}{(4\pi)^{\beta}} \int_0^{\infty} |f(t)|^2 t^{-\beta} dt < \infty$, Duren stated an analogy of the Paley–Wiener theorem for Bergman space.

Theorem A([10]) For each $\alpha > -1$, the space $A^2_{\alpha}(\mathbb{C}^+)$ is isometrically isomorphic under the Fourier transform to the space $L^2_{\alpha+1}(\mathbb{R}^+)$. More precisely, $F \in A^2_{\alpha}(\mathbb{C}^+)$ if and only if it is the Fourier transform $F(z) = \int_0^\infty f(t)e^{2\pi iz \cdot t}dt$ of some function $f \in L^2_{\alpha+1}(\mathbb{R}^+)$, in which case $\|F\|_{A^2_{\alpha}(\mathbb{C}^+)} = \|f\|_{L^2_{\alpha+1}(\mathbb{R}^+)}$. Based on Theorem 1, letting s = 1, p = 2, $\psi(y) = -\frac{\alpha}{4\pi} \log |y|$ and B be a regular open convex cone Γ , we establish Corollary 2, which can be regarded as a higher dimensional and tube domain generalization of Theorem A.

COROLLARY 2. For each $\alpha > -1$, $F \in A^2_{\alpha}(T_{\Gamma})$ if and only if there exists $f(t) \in L^2_{\alpha+1}(\Gamma^*)$ such that

$$F(z) = \int_{\Gamma^*} f(t) e^{2\pi i z \cdot t} dt$$

holds for $z \in T_{\Gamma}$ and $||F||_{A^2_{\alpha}(T_{\Gamma})} = ||f||_{L^2_{\alpha+1}(\Gamma^*)}$.

Proof. By restricting the base B to be a regular open convex cone Γ and letting $\psi(y) = \psi_{\alpha}(y) = -\frac{\alpha}{4\pi} \log |y|, F \in A^2_{\alpha}(T_{\Gamma})$ is also an element of $A^{2,1}(\Gamma, \psi_{\alpha})$. Applying Theorem 1 to such an F, we can show that there exists f(t) satisfying (5) such that $F(z) = \int_{\mathbb{R}^n} f(t)e^{2\pi i z \cdot t} dt$ and $\operatorname{supp} f \subset U_1(\Gamma, \psi_{\alpha})$. Based on (6), we have

$$R_{\psi_{\alpha}} = \lim_{y \in \Gamma, y \to \infty} \frac{\psi_{\alpha}(y)}{|y|} = 0.$$

Thus, together with Lemma 2, the supporter of f(t) is contained in Γ^* and $F(z) = \int_{\Gamma^*} f(t)e^{2\pi i z \cdot t} dt$. Moreover, $\int_{\Gamma} \int_{\Gamma^*} |f(t)|^2 e^{-4\pi (y \cdot t + \psi_{\alpha}(y))} dt dy \leq ||F||_{A^{2,1}(\Gamma,\psi_{\alpha})}$. Thus,

$$\int_{\Gamma} \int_{\Gamma^*} |f(t)|^2 e^{-4\pi (y \cdot t + \psi_{\alpha}(y))} dt dy = \int_{\Gamma^*} \int_{\Gamma} |f(t)|^2 e^{-4\pi y \cdot t} y^{\alpha} dy dt = \int_{\Gamma^*} |f(t)|^2 \frac{\Gamma(\alpha)}{(4\pi t)^{\alpha+1}} dt,$$

which shows $f \in L^2_{\alpha+1}(\Gamma^*)$. And Plancherel's Theorem assures that $||F||_{A^2_{\alpha}(T_{\Gamma})} = ||f||_{L^2_{\alpha+1}(\Gamma^*)}$.

Conversely, note that $F(z) = \int_{\Gamma^*} f(t)e^{2\pi it \cdot z} dt$. For $f(t) \in L^2_{\alpha+1}(\Gamma^*)$, Plancherel's theorem implies that

$$\int_{\mathbb{R}^n} |F(x+iy)|^2 dx = \int_{\Gamma^*} e^{-4\pi y \cdot t} |f(t)|^2 dt,$$
$$\int_{\Gamma} \int_{\mathbb{R}^n} |F(x+iy)|^2 e^{-4\pi \psi_{\alpha}(y)} dx dy = \int_{\Gamma} \int_{\Gamma^*} |f(t)|^2 e^{-4\pi (y \cdot t + \psi_{\alpha}(y))} dt dy < \infty,$$

in which $\psi_{\alpha}(y) = -\frac{\alpha}{4\pi} \log |y|$. Hence, $F(z) \in A^{2,1}(\Gamma, \psi_{\alpha}) = A^2_{\alpha}(T_{\Gamma})$. The proof is complete.

By restricting the base B to be a regular open convex cone Γ , we establish the following weighted version of the edge-of-the-wedge theorem (see [2]) as an application of Theorem 1.

THEOREM 4. Assume that Γ is a regular open convex cone in \mathbb{R}^n , $\psi_1 \in C(\Gamma)$ and $\psi_2 \in C(-\Gamma)$ satisfy

$$R_{\psi_1} = \lim_{y \in \Gamma, y \to \infty} \frac{\psi_1(y)}{|y|} < \infty$$
(18)

and

$$R_{\psi_2} = \lim_{y \in \Gamma, y \to \infty} \frac{\psi_2(-y)}{|y|} < \infty$$
(19)

respectively. If $1 , <math>0 < s(p-1) \leq 1$, $F_1 \in A^{p,s}(\Gamma, \psi_1)$ and $F_2 \in A^{p,s}(-\Gamma, \psi_2)$, satisfying

$$\lim_{y \to 0} \int_{\mathbb{R}^n} |F_1(x+iy) - F_2(x-iy)|^p dx = 0,$$
(20)

then F_1 and F_2 can be analytically extended to each other and further form an entire function F. Furthermore, there exists a function $f \in L^1(\mathbb{R}^n)$ supported in a bounded convex set K such that $F(z) = \int_K f(t)e^{2\pi i t \cdot z} dt$.

Proof. Theorem 1 implies that there exists a function f_j (j = 1, 2) such that

$$F_j = \int_{\mathbb{R}^n} f_j(t) e^{2\pi i t \cdot z} dt$$

holds, in which the supporter of f_j is contained in $U_{sp}((-1)^{j+1}\Gamma, \psi_j)$ for for 1 . Based $on lemma 2, <math>\operatorname{supp} f_j \subset (-1)^{j+1}\Gamma^* + \overline{D(0, R_{\psi_j})}$. By the Hausdorff-Young inequality,

$$\left(\int_{\mathbb{R}^n} |f_1(t)e^{2\pi y \cdot t} - f_2(t)e^{-2\pi y \cdot t}|^q dt\right)^{\frac{1}{q}} \le \left(\int_{\mathbb{R}^n} |F_1(x+iy) - F_2(x-iy)|^p dt\right)^{\frac{1}{p}}.$$

Then it follows from Fatou's lemma and (20) that $||f_1 - f_2||_{L^q(\mathbb{R}^n)} = 0$. Thus, $f_1 = f_2$ almost everywhere on \mathbb{R}^n . Let $f_1(t) = f_2(t) = f(t)$, and $R = \max\{R_{\psi_1}, R_{\psi_2}\}$, then $\operatorname{supp} f \subset K \subset (\Gamma^* + \overline{D(0, R)}) \bigcap (-\Gamma^* + \overline{D(0, R)})$. Thus, K is a bounded convex set. Consequently, $F(z) = \int_K e^{2\pi i z \cdot t} f(t) dt$ is an entire function, where $F(z) = F_1(z)$ for $z \in T_{\Gamma}$ and $F(z) = F_2(z)$ for $z \in T_{-\Gamma}$.

Similarly, we can prove the weighted version of the edge-of-the-wedge theorem for p > 2.

THEOREM 5. Suppose that Γ is a regular open convex cone in $\mathbb{R}^n, \psi_1 \in C(\Gamma)$ and $\psi_2 \in C(-\Gamma)$ satisfy (18) and (19) respectively. If $F_1 \in A^{p,s}(\Gamma, \psi_1)$ and $F_2 \in A^{p,s}(-\Gamma, \psi_2)$, where p > 2, satisfying

$$\lim_{y \in \Gamma, y \to 0} \int_{\mathbb{R}^n} |F_1(x+iy) - F_2(x-iy)|^2 dx = 0,$$
(21)

then F_1 and F_2 can be analytically extended to each other and further form an entire function F. Furthermore, there exists a measurable function f(t) supported in a bounded convex set K such that $F(z) = \int_K f(t)e^{2\pi i t \cdot z} dt$.

Proof. For $F_j \in A^{p,s}((-1)^{j+1}\Gamma, \psi_j)$ and $\frac{1}{p} + \frac{1}{q} = 1$, exists a measurable function f_j such that $F_j = \int_{\mathbb{R}^n} f_j(t) e^{2\pi i t \cdot z} dt$ and $\operatorname{supp} f_j \subset U_{sp}((-1)^{j+1}\Gamma, \psi_j)$, where j = 1, 2. It then follows from

Lemma 2 that $\operatorname{supp} f_j \subset (-1)^{j+1} \Gamma^* + \overline{D(0, R_{\psi_j})}$. Plancherel's Theorem implies that

$$\left(\int_{\mathbb{R}^n} |f_1(t)e^{2\pi y \cdot t} - f_2(t)e^{-2\pi y \cdot t}|^2 dt\right)^{\frac{1}{2}} = \left(\int_{\mathbb{R}^n} |F_1(x+iy) - F_2(x-iy)|^2 dx\right)^{\frac{1}{2}}.$$

Then based on (21) and Fatou's Lemma, $||f_1 - f_2||_{L^2(\mathbb{R}^n)} = 0$, which means $f_1 = f_2$ almost everywhere on \mathbb{R}^n . Let $f_1(t) = f_2(t) = f(t)$ and $R = \max\{R_{\psi_1}, R_{\psi_2}\}$, then $\operatorname{supp} f(t) \subset K = (\Gamma^* + \overline{D(0, R)}) \bigcap (-\Gamma^* + \overline{D(0, R)})$. Thus, K is a bounded convex set. As a result, $F(z) = \int_K e^{2\pi i z \cdot t} f(t) dt$ is an entire function, where $F(z) = F_1(z)$ for $z \in T_{\Gamma}$ and $F(z) = F_2(z)$ for $z \in T_{-\Gamma}$.

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